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Vaporization and Phase Equilibria of Simulated Radionuclides

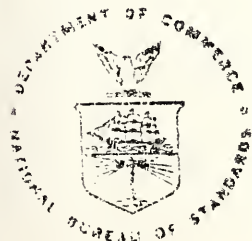
E.R. Plante and J.W. Hastie

U.S. DEPARTMENT OF COMMERCE
National Bureau of Standards
Institute for Materials Science and Engineering
High Temperature Chemistry
Ceramics Division
Gaithersburg, MD 20899

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January 1986

Interim Report



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NATIONAL BUREAU OF STANDARDS

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U.S. DEPARTMENT OF COMMERCE, Malcolm Baldrige, *Secretary*
NATIONAL BUREAU OF STANDARDS, Ernest Ambler, *Director*

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Vaporization and Phase Equilibria of Simulated Radionuclides

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ABSTRACT

Undesirable losses of radionuclides occur by vaporization during processing of nuclear waste glass. This study examines the phase and chemical transformations that occur during the initial steps of waste incorporation into a borosilicate glass host. Specialized mass spectrometric and thermogravimetric techniques were used to monitor the vapor transport species over a range of temperature and composition. Significant vapor transport of Cs was found and its enhancement by the presence of carbon, halogen, and water vapor was also noted.

Key words: Cesium, mass spectrometry, nuclear waste, thermodynamics, vaporization, CsOH , Cs_2O , Cs_2CO_3 , CsO_2CH .

1. Introduction

Industrial plans for nuclear waste storage call for the waste to be incorporated in borosilicate glass. A large fraction of the nuclear waste, known as "sludge", consists of both insoluble and hydrous oxides which have been reduced by formic acid. This sludge contains about 95% of the nuclear waste which, according to present plans, will be disposed of by vitrification in borosilicate glass at temperatures of 850-1150°C [Bickford, 1985]. Typical sludge and melter feed sludge-glass compositions are given in table 1. Plant experience suggests that losses of radionuclides, and particularly cesium,

Table 1

Chemical Compositions of Sludge and Melter Feed Simulations
[Jantzen et al. 1982]

Chemical	Sludge (wt%)	Melter feed (wt%)	Chemical	Sludge (wt%)	Melter feed (wt%)
SiO ₂		42.648	La ₂ O ₃	0.487	0.182
B ₂ O ₃		6.272	ZrO ₂	0.446	0.166
Na ₂ O		8.153	CrCl ₃ ·6H ₂ O	1.095	0.408
Li ₂ O		4.390	Ag	0.078	0.029
MgO		0.627	Cu ₂ O	0.083	0.031
ZrO ₂		0.627	CoCl ₂ ·6H ₂ O	0.012	0.004
Fe ₂ O ₃	24.769	9.234	Zn	0.165	0.062
Mn(COOH) ₂	12.321	4.594	Mg(NO ₃)·6H ₂ O	1.792	0.668
UO ₂ (NO ₃) ₂ ·6H ₂ O	9.040	3.370	Na ₃ PO ₄ ·12H ₂ O		0.020
Al(OH) ₃	15.567	5.804	NaF		0.030
CaCO ₃	3.437	1.281	Cs ₂ O		0.028
CaSO ₄	0.171	0.064	Coal	0.063	0.023
PO ₄	0.094	0.035	Zeolite ^b	5.654	2.108
CaF ₂	0.080	0.030	Na ₂ SO ₄	0.148	0.055
NiO	2.004	0.747	KOH	0.249	0.093
NaCl	0.273	0.102	Sr(NO ₃) ₂	0.174	0.065
HgO	1.804 ^a	0.673 ^a	HCOOH(90%)	11.5 ^c	4.03 ^c
NaI	0.008	0.003			
RuO ₂	0.241	0.090	SiO ₂	0.464 ^d	0.883 ^d
Na ₂ SiO ₃ ·9H ₂ O	18.063	6.735	Al ₂ O ₃	0.108 ^d	0.040 ^d
ThO ₂	0.637	0.237	B ₂ O ₃	0.000 ^d	0.090 ^d
BaO	0.227	0.085	Alkali	0.011 ^d	0.560 ^d
Ce(SO ₄) ₂	0.636	0.237	Alkali earths	0.052 ^d	0.035 ^d
PbSO ₄	0.345	0.129	Fe	0.234 ^d	0.087 ^d
			Other transition	0.181 ^d	0.073 ^d
			metals	0.234 ^d	0.087 ^d
			Rare earths	0.030 ^d	0.001 ^d
			U, Th	0.020 ^d	0.008

^aNot added to simulation

^bLinde AW500

^cComposition expressed as cm³/100 g

^dComposition expressed as mol/100 g

occur by vaporization. In a previous study [Hastie et al., 1983], we characterized these losses from simulated (non-radioactive) melter feed compositions. The present study focuses on the earlier processing stage of sludge heating and mixing with the glass host.

Because of the highly complex chemical and phase characteristics of sludge, it was necessary, in the present study, to examine key sludge components as well as the sludge itself. The component data provide the characteristic mass spectral ionization patterns and cross sections necessary to the analysis of complex mixtures. Also, the literature thermodynamic function and vapor pressure data for the reference state components is largely incomplete or inaccurate. These data are essential to the interpretation of experimental data and to the development of thermodynamic mixing models for complex mixtures such as the melter feed and sludge systems considered here. Emphasis has been given to determination of the vaporization characteristics of key cesium-containing compounds (non-active isotopes) and the vaporization behavior of synthetic sludge and sludge incorporated in the borosilicate glass.

Note in table 1, that the chemical composition of the typical melter-feed simulation (ie, non-radioactive isotopes) differs from that of our previous study [Hastie et al., 1983]. The major oxide components differ typically by a few wt % in each case. Also, the present composition contains salts such as sulfates, nitrates, and halides. These sulfate and nitrate forms should react at relatively low temperatures to form oxides which are then incorporated into the glass. However, the halides may be more difficult to incorporate.

2. Experimental Techniques

2.1 Knudsen Gravimetric Analysis (KGA)

With this classical technique, vaporizing species effuse from a container-cell through a small orifice (~ 1mm diameter) into a high vacuum. The equation

$$P = \left(\frac{m}{CA t} \right) \left(\frac{2\pi RT}{M} \right)^{1/2} \quad (2.1)$$

relates the vapor pressure (P) to the rate at which mass (m) is transferred through an orifice (area A); where t is time, R the universal gas constant, T the absolute temperature, and M the molecular weight of the effusing species. The factor C corrects for the transmission of the vapor through a non-ideal orifice (i.e., not infinitely thin). This relationship is based on the condition of molecular effusion where the molecular mean free path is larger than the orifice diameter. Generally, this places an upper pressure limit on the validity of the method at about 10^{-4} atm. This approach can be used to determine precise pressures but if more than a single vapor species is produced, and these are unrelated by a chemical constraint, then additional information is needed to partition the vapor species between the two or more competing reactions.

For the present study, a Mettler thermogravimetric analyzer was used.

2.2 Knudsen Mass Spectrometry (KMS)

This method combines the classical Knudsen technique, described above, with a mass spectrometric detector. The molecular beam emitted from the cell is passed through a rotating disk which modulates the beam. This procedure

allows for phase sensitive detection which distinguishes beam gases from background. The technique is especially useful for permanent gases which otherwise would lead to untenable signal/noise ratios. Ion currents, produced by electron impact ionization, are converted to pressure using the expression

$$P = kI_iT, \quad (2.2)$$

where the partial pressure of species i is related to the observed ion signal, I_i , at temperature T , by a system constant, k , which contains ionization cross section, mass filter transmission, and geometric factors. This constant is usually determined by methods which depend on ion current data, integrated with respect to time and temperature, together with gravimetric data, as discussed in detail elsewhere [Hastie et al., 1983]. For the sludge-containing systems, the vapor composition was very complex and conversion of mass spectral ion intensities to partial pressures was accomplished using a constant average value of k for all species. Errors of up to a factor of two may arise with this approximation.

With the gas-inlet KMS system, the Knudsen cell is modified to allow the introduction of small quantities of gas into the cell. Hence, it is possible to study the interaction of the added gas with the condensed phase material in the cell.

2.3 Transpiration Mass Spectrometry (TMS)

The TMS technique uses a platinum transpiration cell, typically operating in the 0.1-1 atm total pressure range. A small orifice (0.08 mm diameter), at the downstream end of the cell, serves as a sonic nozzle which produces a "frozen equilibrium" molecular beam. This beam is then directed toward a quadrupole mass spectrometer system in a specially designed high pumping

speed vacuum system, as described earlier [Hastie, et al., 1983]. The differentially pumped molecular beam is modulated to isolate molecular beam components from scattered gas.

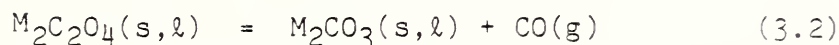
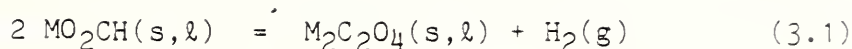
An advantage of the high pressure inlet system of TMS is that it decreases bubbling phenomena which can be a problem with the KMS method. Another advantage is that the much higher pressures present, and the relatively small orifice, lead to an increased residence time (seconds versus milliseconds for the KMS system) and hence a closer approach to thermodynamic equilibrium. The TMS technique is also the only method which provides mass analysis of vaporization products above the pressure range allowed in Knudsen effusion measurements. It should be recognized, however, that because of the low beam temperature, the characteristic electron impact mass spectrum can differ significantly from that found in KMS measurements [Bonnell and Hastie, 1979]. For certain alkali salts this leads to difficulty in detecting parent molecular ions.

Vapor pressures can be related to ion currents using the known carrier gas pressure as a reference. Alternatively, ion current integrations can be used, in combination with measured mass losses, to determine the instrument sensitivity, as discussed earlier [Hastie et al., 1983].

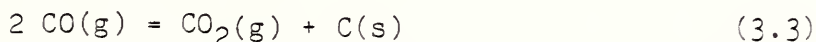
3. Vaporization Studies of Selected Sub-Systems of Nuclear Waste

3.1 Cesium formate

Meisel et al. [1975] have studied the decomposition of alkali metal (M) formates using thermogravimetric analysis (TGA). The results indicate the reactions:



and these may proceed stepwise or simultaneously. In addition, the direct observation of a small amount of carbon in the reaction product suggests that another reaction:



takes place catalytically during the oxalate decomposition step (3.2). Meisel et al. [1975] found evidence for the oxalate intermediate step for sodium and potassium formate but not cesium which proceeded directly to the carbonate and was independent of heating rate, atmosphere or container material (glass or platinum). However, our TGA experiments in vacuum indicated that the two reactions can be separated if the temperature is changed stepwise. At 210°C, reaction (3.1) proceeded to completion, although at a slow rate (time scale ~ 2 hrs). On raising the temperature to 325°C no further weight loss occurred. At 425°C, reaction (3.2) commenced and was driven to completion at 475°C.

Both the KMS and TMS mass spectrometer results clearly confirmed reaction (3.1), with the observation of H_2 but not CO (or CO_2) at low temperatures. It is noteworthy that Cs^+ was present in the mass spectrum during this stage and the signal disappeared as the H_2 emission fell at the end of reaction (3.1). This could be due to an here-to-fore undocumented vaporization of CsO_2CH as a stable vapor species. Alternatively, at this low a temperature (190°C), both Cs(l) and, to a much lesser degree, CsOH(l) are sufficiently volatile to be transported and either might be formed by reaction with H_2 . In fact, both the Cs and H_2 mass spectrometric signals decreased with time together, which indicates that Cs is most likely a reaction product of H_2 with the sample. The TGA results also indicated the

same coupled loss of Cs and H₂ but at 228°C. The higher temperature required in this case was due to the lower sensitivity of the TGA method, relative to KMS or TMS. Analysis of the TGA Knudsen effusion data indicated that about one percent of the initial Cs content of the sample was lost to the vapor phase during H₂ release. No further weight loss was observed on raising the temperature through 325°C.

During the second step (3.2) CO emission was evident, and the final residue was predominately Cs₂CO₃(c). TGA analysis indicated that reaction (3.2) began at 425°C and was driven to completion by 475°C. The product contained a small amount of elemental carbon which is a strong O₂ scavenger. With the O₂ partial pressure lower by several orders of magnitude (below the ~ 10⁻⁸ atm detection limit of KMS), the environment is reducing and the resulting partial pressure of Cs over formate-derived Cs₂CO₃ is roughly an order of magnitude greater than that observed from reagent-grade Cs₂CO₃, as discussed below (section 3.4).

TMS measurements carried out at a total pressure of 0.3 atm (N₂ carrier gas) also showed the loss of Cs associated with H₂ evolution. With more sample than was used in the Knudsen effusion experiments, and with higher heating rates, evolution of H₂, CO, and Cs was observed at a total carrier gas pressure of 1 atm and at temperatures near 500°C before exhaustion of H₂. At temperatures in the 500-700°C region, appreciable pressures of CO and Cs, as well as a significant CO₂ signal, were observed. The high Cs signals most likely resulted from reduction of the formate-derived Cs by the C(s) formed during the decomposition process. Further discussion of this process is given below in section 3.4.

3.2 Cesium Hydroxide

The JANAF thermochemical tables [JANAF, 1971] provide derived heat of formation data for CsOH which should be considered an approximation. This derivation is based on: atomic absorption spectroscopic observation of Cs-depletion in flames, electron impact mass spectrometric data, and bond energy analogy with other alkali hydroxides. Besides the monomer, the dimer(CsOH)₂ has also been observed by mass spectrometry. A derivation of the dimer heat of formation was made by JANAF [1971] using the equilibrium data of Schoonmaker and Porter [1959] for the interaction of KOH and CsOH to form the respective dimers. From these results, the difference in the heat of dimerization of KOH and CsOH was obtained. The heat of formation of the dimer is thus based on that for (KOH)₂ and this measured difference. It should be noted that the (KOH)₂ heat of formation given by JANAF [1971, 1975] has been revised recently, as discussed elsewhere [Hastie, et al. 1984] and the JANAF (CsOH)₂ data should therefore be revised.

Accurate vapor pressure measurements of CsOH are difficult to make because of interaction with the cell material, its tendency to creep from the cell and vaporize from external surfaces, and the possibility of reaction with CO₂ after preparation and handling such that the activity of CsOH(l) is reduced.

In the present study, vapor pressure data were obtained using both the TGA and KMS techniques. The pressures obtained from the TGA Knudsen effusion experiment, as given in figure 1, were calculated assuming that CsOH(g) was the only vapor species. Based on the KMS results, and related discussion given below, the true molecular weight lies between that of CsOH and (CsOH)₂. The pressure scale of the figure 1 TGA data is probably about 10 percent too

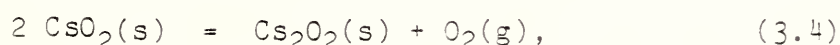
high, which is similar to the experimental uncertainty. A second law analysis of the data gives a heat of vaporization at 298 K of 27.8 kcal/mol while the third law value is 32.6 kcal/mol. The JANAF [1971] value is 35.0 kcal/mol.

KMS experiments using the CsOH sample from the TGA experiment indicated Cs^+ , CsOH^+ and Cs_2OH^+ mass spectral ions. At a nominal 30eV ionizing electron energy, the following relative ion intensities were found: Cs^+ (100), CsOH^+ (14), and Cs_2OH^+ (10). Thus the dimer is present at approximately 10 percent of the monomer partial pressure, in marked contrast to the literature value as evaluated by JANAF [1971], eg. see figure 1. This difference could be due, in part, to the following observation. Increasing the temperature after depletion of the CsOH from the Knudsen cell showed that the sample contained Cs_2CO_3 as indicated by the appearance of CO_2^+ and O_2^+ . The presence of a dissolved impurity, such as carbonate, would decrease the CsOH thermodynamic activity. Since the CsOH pressure is proportional to the activity, while the dimer pressure is proportional to its square, impurities will have a greater effect on the dimer pressure. However, the original sample would have to contain at least 70 percent dissolved Cs_2CO_3 to account for the difference between our KMS data and JANAF's. Such an impurity level, and degree of carbonate solubility, is highly unlikely. Also, the high pressure found by KGA (in fig. 1) suggests a CsOH activity near unity. Thus we conclude that the true dimer concentration falls closer to the 10 percent value observed in this work than to the literature value of 80 percent (JANAF, 1971). Based on the present mass spectrometric observations, the dimer concentration is estimated to be 25 ± 10 percent which would indicate an average molecular weight for the KGA measurements (fig. 1) of 187 rather than the assumed 150 amu. This dimer concentration is consistent with the trends observed for the other alkali hydroxide systems, eg. see Hastie et. al [1984].

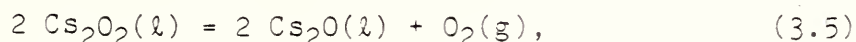
Although there is considerable uncertainty with respect to the relative importance of the monomer and dimer, both in this and the earlier literature work, it is clear that the high volatility of the hydroxide may lead to significant transport of Cs in the presence of a high H₂O partial pressure in the glass-sludge reactor atmosphere. Future measurements should be made over solid CsOH, using higher sensitivity mass spectrometry, to avoid the liquid creep and reaction problems encountered.

3.3 Cesium Oxide

Lamoreaux and Hildenbrand [1984] and Glushko et al. [1978-1982] have recently reviewed the literature and recommended thermodynamic functions for the more thermally stable cesium oxide compounds, as summarized in table 2. The agreement between the two data sets is reasonably good except for Cs₂O₂(c). Lamoreaux and Hildenbrand [1984] used reaction enthalpy data from Berardinelli and Kraus [1974] for



together with the heat of formation of CsO₂(s). Glushko et al. [1978-1982] used data from Centnerzwer and Blumenthal [1933] based on the reaction,



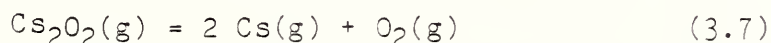
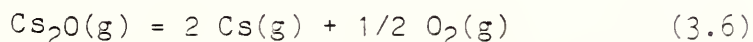
together with the heat of formation of Cs₂O(l). However, the data of Centnerzwer and Blumenthal [1933] is probably not reliable because it was taken at temperatures where only a single liquid phase exists, contrary to reaction (3.5). Hence the value given by Lamoreaux and Hildenbrand [1984] is preferred.

Table 2

Heats of Formation and Entropies of Cs-O Compounds at 298 K

Compound	Lamoreaux and Hildenbrand [1984]		Glushko et al. [1978-1982]	
	ΔH_f kcal mol ⁻¹	S° cal mol ⁻¹ K ⁻¹	ΔH_f kcal mol ⁻¹	S° cal mol ⁻¹ K ⁻¹
Cs(c)	0	20.37 ± 0.10	0	20.37
Cs(g)	18.32 ± 0.50	41.95 ± 0.00	18.28	41.94
CsO(g)	9.73 ± 1.0	59.54 ± 0.60	8.98	59.37
CsO ₂ (c)	-68.40 ± 0.50	33.40 ± 1.5	-68.38	33.94
CsO ₂ (g)	2.00 ± 5	67.78 ± 1.5	-	-
Cs ₂ (g)	27.30 ± 1.0	67.86 ± 0.20	26.15	68.01
Cs ₂ O(c)	-82.69 ± 0.30	35.09 ± 0.10	-82.79	35.11
Cs ₂ O(g)	-37.96 ± 6.0	75.97 ± 2.0	-34.14	77.44
Cs ₂ O ₂ (c)	-119.00 ± 2.5	34.50 ± 1.5	-105.16	43.02
Cs ₂ O ₂ (g)	-60.00 ± 6.0	81.49 ± 2.0	-59.05	81.46

The derived heats of formation for $\text{Cs}_2\text{O}_2(\text{g})$ and $\text{Cs}_2\text{O}(\text{g})$ are based on two separate mass spectrometric vaporization studies by Gusarov et al. [1967a] and Norman and Staley [1966]. Gusarov et al. [1967a] used $\text{CsO}_2(\text{s})$ as a starting material which was then thermally decomposed in the Knudsen cell to form $\text{Cs}_2\text{O}_2(\text{s})$. Following this thermal decomposition step and the evaporation of CsOH impurity, mass spectral ion intensity data (I^+) were obtained for O_2^+ , Cs^+ , Cs_2O^+ , and Cs_2O_2^+ . Their plots of $\log \text{I}^+T$ vs $1/T$ yielded straight lines, as required to a first approximation by the second law of thermodynamics. However, the slopes apparently did not correspond to meaningful enthalpies of sublimation or vaporization because of sample interaction with the platinum effusion cell. These data were therefore used to calculate heats for the all-gas reactions

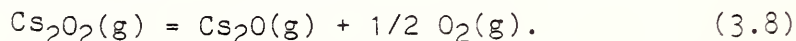


which, in principal, are independent of the composition of the condensed phase.

Other data given by Gusarov et al. [1967a] indicate oxygen pressures at about ten times higher than those required to maintain $\text{Cs}_2\text{O}_2(\text{s})$ at the stoichiometric composition. In addition, the mass spectral ion with the highest intensity was Cs^+ which was formed by ionization of $\text{Cs}(\text{g})$. Also, plots of the ion ratios Cs^+/O_2^+ and $\text{Cs}^+/\text{Cs}_2\text{O}_2^+$ as a function of vaporization time, shown in figure 2, indicate a shift in composition towards a more Cs-rich sample.

The second mass spectrometric study, by Norman and Staley [1966], provides additional data on the heat of formation of $\text{Cs}_2\text{O}(\text{g})$ and $\text{Cs}_2\text{O}_2(\text{g})$. However, those authors did not detect a significant Cs^+ signal over $\text{Cs}_2\text{O}(\text{c})$, in contrast to the measurements of Gusarov et al. [1967a]. This observation indicates considerable uncertainty concerning the composition of the

condensed phase in most of these measurements. Norman and Staley's measurements were made by passing O_2 through a reentrant tube into the Knudsen cell to establish a relatively high O_2 pressure (perhaps 10^{-4} atm). They monitored reaction (3.6) in addition to



For reaction (3.8), the sample composition was believed to be between $Cs_2O_2(c)$ and $Cs_2O(c)$ but not very close to either, while for reaction (3.6), cesium silicate was used as the source of Cs. The third law data obtained are similar to the values recommended by Lamoreaux and Hildenbrand [1984].

In this study, KMS vaporization experiments were conducted on $Cs_2O(c)$ using a sample obtained from Alfa in the as-received condition. The sample consisted of large purplish-greenish chunks and, to minimize reaction with the ambient atmosphere, a single chunk was inserted into a platinum Knudsen cell and loaded into the mass spectrometer as rapidly as possible. Even with these pre-cautions mass spectral ions (at m/e of 133, 150, and 283 amu), corresponding to the presence of condensed $CsOH(c)$, were identified during the initial heating period. The disappearance of the 150 and 283 amu ion signals with time indicated that $CsOH$ had been purged from the sample. In the initial heating stage the O_2^+/Cs^+ ratio was unusually high, indicating that the original $Cs_2O(c)$ sample contained O_2 in excess of the stoichiometric amount.

Following this initial heating period, data were obtained for Cs^+ , O_2^+ , and Cs_2O^+ . From the $\log I^+T$ vs $1/T$ results, the $\Delta H(T)$ values for Cs, O_2 and Cs_2O were 39.1, 39.0, and 49.5 kcal mol $^{-1}$, respectively. When treated as though reaction (3.6) took place, a second law heat of 53.7 kcal mol $^{-1}$ was obtained. The equilibrium constant data for this reaction are given in figure 3, together with the recommended data of Lamoreaux and Hildenbrand

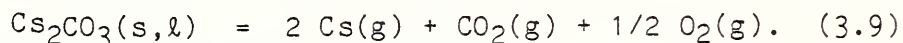
[1984] for comparison. Using an estimated value of the mass spectrometer constant (k), a third law reaction enthalpy of $64.8 \text{ kcal mol}^{-1}$ was obtained. The third law value is fairly insensitive to the accuracy of the mass spectrometer constant and is likely to be more reliable than the second law value. From the third law reaction enthalpy, we obtain $\Delta H_f(298) \text{ Cs}_2\text{O(g)}$ equal to $-28.2 \text{ kcal mol}^{-1}$. This value is about 10 and 6 kcal mol^{-1} less negative than the values recommended by Lamoreaux and Hildenbrand [1984] and Glushko et al. [1978-1982], respectively. The close agreement between the vaporization enthalpies for Cs^+ and O_2^+ and the relative ion intensities indicate that the parents of these ions (Cs and O_2) are most likely both formed by thermal dissociation of $\text{Cs}_2\text{O(g)}$ (i.e. reaction 3.6).

Measurements of the vaporization enthalpy for $\text{Cs}_2\text{O(c)}$ have also been reported by Nicolosi, Tang, and Munkelwitz [1979]. They also used a phase-locked mass spectrometer system. Although Cs^+ was the most abundant ion they assumed it was formed as an electron impact fragment of $\text{Cs}_2\text{O(g)}$. From a plot of $\log I^+T$ vs $1/T$ for Cs_2O^+ they obtained a ΔH_v ($\sim 1000 \text{ K}$) value of $62.9 \text{ kcal mol}^{-1}$. The expected value, based on Lamoreaux and Hildenbrands' [1984] analysis, is $35.7 \text{ kcal mol}^{-1}$. Using a similar assumption for the source of Cs^+ , our data would yield $49.5 \text{ kcal mol}^{-1}$.

We also attempted to reproduce the results of Norman and Staley [1966] by passing $\text{O}_2(\text{g})$ over the $\text{Cs}_2\text{O(c)}$ phase. This approach should have allowed us to determine an equilibrium constant for reaction (3.8). However, this was not possible as we were unable to observe Cs_2O_2^+ . Evidence for O_2 absorption was provided by an observed decrease in the Cs^+ and Cs_2O^+ signals. Also, on increasing the sample temperature, there was an increase in the O_2^+ signal. This apparently means that the oxygen pressure over $\text{Cs}_2\text{O}_{1+x}$ is somewhere in the 10^{-4} atm range at 900 K .

3.4 Cesium Carbonate

Vaporization measurements were carried out using both the TGA and KMS systems. Samples included pure Cs_2CO_3 and Cs_2CO_3 derived from the thermal decomposition of CsO_2CH . As with other alkali metal carbonates, the most important vaporization process is



Gusarov et al. [1967b] have demonstrated that a negligibly small fraction of the saturated vapor consists of $\text{Cs}_2\text{CO}_3(\text{g})$ and $\text{Cs}_2\text{O}(\text{g})$. These minor constituents were also observed in our KMS measurements but will not be of further interest here.

Data from the KGA Knudsen experiments, using both pure carbonate and formate-derived carbonate, are shown in figure 4. Accepting that the decomposition reaction (3.9) is dominant, one can show that the Cs pressure is 1.28 times the effective pressure calculated using the molecular weight of Cs_2CO_3 . The calculated pressures for the formate-derived carbonate are 10 to 100 times greater than the pressures measured using the pure carbonate. This is caused by the carbon residue formed during the decomposition of the formate, as noted above in section 3.1. Note in figure 4 that the Cs pressures over pure carbonate agree within a factor of 3 with those calculated from the thermodynamic functions given by Glushko et al [1978-1982] while the pressures for the formate-derived carbonate are in reasonable agreement with the pressures calculated for carbon-reduced carbonate.

In the KMS experiments on the formate-derived carbonate, the dominant ions were Cs^+ , CO_2^+ and CO^+ . The O_2^+ signal was below the detection limit of the mass spectrometer and this may be contrasted with the pure carbonate case

where O_2^+ was readily observable. This result also confirms the reduction process of the formate-derived carbonate. Although the amount of carbon formed in the decomposition of the formate is believed to be small, CO^+ signals were persistent throughout the sample vaporization. However, all signals decreased with time, showing that kinetic factors were causing a deviation from equilibrium which is typical of reactions involving two condensed phases.

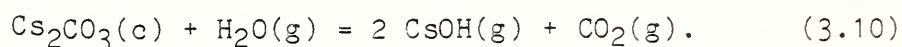
The KMS Cs pressure results obtained over pure Cs_2CO_3 are shown in figure 5. For these measurements the mass spectrometer calibration constant was determined from the measured weight loss and the mass spectrometric data expressed as $[I \cdot t \cdot T^{1/2}]$. The measured Cs pressure is roughly a factor of 10 greater than the Cs pressure calculated from the thermodynamic data tabulated by Glushko et al. [1978-1982]. A possible explanation for this difference is that a shift in sample composition takes place under the vacuum vaporization conditions. Reisman [1958] has reported that alkali carbonates, including Cs_2CO_3 , melt at lower temperatures in the absence of CO_2 than they do in a CO_2 atmosphere. This effect could be caused by partial decomposition to form an oxide-carbonate solution in the absence of a relatively high CO_2 pressure.

3.5 Cesium Carbonate - Water Vapor

The relatively high partial pressure of H_2O in the off-gas for the nuclear waste process leads to the possibility of water interaction with Cs salts dissolved in the borosilicate glass. In our previous work [Hastie et al., 1983] with Cs-containing borosilicate glass we were unable to observe any enhancement in Cs transport in the presence of water vapor. This lack of reactivity was attributed to the low thermodynamic activity of Cs in the glass

which was probably bound as CsBO_2 . On the other hand, Carpenter et al. [1983] did observe an H_2O enhancement of Cs transport over Synroc ($\text{Ba}_{0.9}\text{Cs}_{0.1}\text{Al}_2\text{Ti}_6\text{O}_{16}\text{TiO}_2$), but in this case the Cs is undoubtedly present at a higher activity.

The conversion of CsO_2CH to carbonate may provide a potential source for water reaction with formate-derived carbonate. Therefore measurements were made on the effect of H_2O on Cs_2CO_3 evaporation using the gas-inlet KMS system. Observation of the 150 amu species, CsOH^+ , and an enhanced CO_2^+ signal indicated the reaction



Data obtained for this process at 1010 K are plotted in figure 6. The curve with zero slope shows that the reaction equilibrium constant, expressed as an ion intensity quotient, is independent of the ion intensity (and hence pressure) of H_2O . This result indicates establishment of equilibrium or, at the least, a steady state with respect to reaction (3.10). The second curve has a slope of 0.97 which is in excellent agreement with the theoretical slope of 1.0 for reaction (3.10).

Absolute equilibrium constant data were calculated from the results given in figure 6 using the mass spectrometer constant (k). Figure 7 compares the experimental data with the curve calculated from the tabulation of Glushko et al. [1978-1982]. The experimental curve is about three orders of magnitude lower than the calculated curve but this is not unusual with Knudsen effusion gas-solid experiments where the reactant gas residence time is short (usually milliseconds). At least part of this difference can be attributed to the considerable uncertainty in the literature Gibbs energy data for the CsOH species (see section 3.2) used to generate the calculated curve of figure 7. These results show that cesium vapor transport can be

enhanced significantly by water vapor, in accordance with reaction (3.10). However, this enhancement may be negligible in practice if the Cs activity is reduced by reaction with the borosilicate glass.

3.6 Sludge

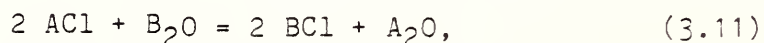
3.6.1 KMS Observations

Experiments were carried out with samples having the composition given in table 1. The samples were contained in a platinum Knudsen cell. Because of the high sensitivity of the KMS technique, decomposition products can generally be observed at lower temperatures than for the TMS method or for actual reactor plant conditions. However, to maintain pressures in the Knudsen cell within the applicable limits ($P < 10^{-4}$ atm), a relatively low heating rate (~ 1 °C/min) was necessary.

Decomposition of formates in the sludge was observed at 190-200°C. Prominent mass peaks occurred at 2, 18, 28, 30, and 44 amu. The mass peaks at 2 and 28 amu are characteristic of formate decomposition. Mass 18 amu is clearly water vapor which originates as residual water or as salt hydrates. The 44 amu signal may be associated with the reaction of CO to form CO₂ and C(s), as well as from the decomposition of residual carbonates. The origin of mass 30 amu is somewhat uncertain. This peak was not observed in the cesium formate mass spectrum but, nevertheless it does appear to be very closely related to formate decomposition in the sludge. The possibility that this peak corresponds to the NO radical was considered but the temperature is too

low for decomposition of nitrate salts. The most likely assignment, is to formaldehyde, which may arise as an impurity in the formic acid used in the reduction process.

At temperatures near 500°C alkali salts began to vaporize from the sludge, as evidenced by the appearance of Na^+ , K^+ , NaCl^+ , and KCl^+ . The mass spectral ion Na_2Cl^+ , resulting from $(\text{NaCl})_2$ was also observed. The presence of the dimer suggests that NaCl(g) is near its saturation pressure and the condensed salt activity is near unity. Near 800°C, the NaF^+ ion was also observed. This species assignment could be distinguished from LiCl^+ (at the same amu) because of the absence of Li^+ in the mass spectrum. At the temperatures involved, all the alkali metal ions result from mass spectral fragmentation of parent halide molecules. It may be noted that the sludge composition, as prepared, contains no KCl suggesting that this salt is formed by exchange reactions with other halides, as we have noted in other studies [Hastie and Plante, 1981]. These condensed phase reactions have the form



where A and B are different alkali (or even alkaline-earth and transition) metals and the oxides may also be combined as silicates, borates, etc.

Also, near 500°C the sulfate salts began to decompose, as indicated by the appearance of SO_2^+ , SO^+ , and S^+ , which are characteristic of the SO_2 mass spectrum. The sludge melted at about 625°C and the signals became erratic at higher temperatures due to nucleation of bubbles formed from gaseous products which migrated through the melt and burst at the surface. The major component of the bubbles appeared to be SO_2 . However, the intensities of other mass peaks were also affected because bulk melt components, with higher

concentrations than at the quiescent surface, can be transported to the surface by bubbling. Near 1000°C, Pb^+ was observed indicating vaporization of Pb(g) from the sludge.

3.6.2 TMS Observations

TMS experiments were carried out using an argon carrier gas. Because of the less stringent pressure requirement for the TMS system, the heating rate was faster ($\sim 10^\circ\text{C/min}$) than for the KMS experiments. Typical species evolution data, obtained for a single heating cycle over the range 150-450°C, are shown in figure 8. The observed signals indicate thermal decomposition of the formates and other relatively unstable sludge components. Note that the curves for H_2O^+ , CO_2^+ , and H_2CO^+ have maxima near 280°C. These correspond to the major signals seen in the KMS experiments near 200°C. The low temperature at which these gases are expelled suggests that they are not strongly bonded in the sludge substrate. CO_2 and H_2O could originate from decomposition of bicarbonates, relatively unstable carbonates such as PbCO_3 , and from salt hydrates. As with the KMS observations, we believe that the H_2CO^+ signal originates from a formaldehyde impurity.

The second group of peaks, with maxima near 380°C, are formed from formate decomposition but the rate of heating was high enough to allow the formate-oxalate and oxalate-carbonate transformations to occur simultaneously. As with the formates, CO_2 is due to the decomposition of CO which is thermodynamically unstable with respect to CO_2 and C(s) in this temperature range. A third CO_2 intensity maximum near 425°C, presumably results from the decomposition of a relatively stable carbonate.

At higher temperatures (than shown in fig. 8) Na^+ , SO_2^+ , SO^+ , and S^+ were observed and at the highest temperatures, following the decay of the SO_2^+ signal, the O_2^+ mass spectral species was observed. The Na^+ signal, which was observed throughout the temperature range, undoubtedly results from mass spectral fragmentation of NaCl(g) at the lower temperatures and from decomposition of $\text{Na}_2\text{SiO}_3\text{(s)}$ or $\text{Na}_2\text{SO}_4\text{(s)}$ at the higher temperatures. Note that K^+ was not observable in the TMS experiments because of interference by the Ar carrier gas at the 40 amu position.

The SO_2 pressure showed relatively little temperature dependence but was gradually depleted from the sample with time. For Ar pressures of 0.2-0.4 atm, no apparent SO_2 bubble formation occurred, in contrast to the KMS experiments. The O_2 pressure showed a temperature dependence similar to that calculated for $\text{Fe}_2\text{O}_3\text{-Fe}_3\text{O}_4$ reduction, indicating further reduction of Fe_2O_3 in the sludge. However, for the off-gas conditions present in the industrial glass reactor, loss of oxygen by reduction of Fe_2O_3 probably would not take place.

3.7 Sludge-Glass Frit Mixture

The sludge-frit or melter-feed mixture, as prepared by SRL [Bickford, 1985], consisted of sludge and a borosilicate glass frit, mixed in the respective weight proportion of 1/1.682, which was fused and quenched. The overall composition of the sludge-frit mixture, as given in table 1, is similar, although not identical, to the synthetic nuclear waste glass studied by us earlier [Hastie, et al., 1983]. Several samples were studied and the melting and quenching performed during preparation of the mixture appears to have removed most of the relatively volatile material associated

with the formate salts. It is also likely that some of the glass-forming salts, such as sulfates, may have undergone partial reaction during the mixture preparation.

3.7.1 KMS Observations

The vaporization behavior of the mixture was quite erratic throughout the temperature range covered of 500-1150°C. It appears that the sludge-frit mixture consists of at least two phases, one containing salt and the other glass-forming components. Such a phase separation would contribute to the erratic vaporization behavior.

In the 550-700°C temperature range, Cs vapor transport was indicated by the appearance of Cs^+ , as shown in figure 9. At such low temperatures, CsCl , CsOH , or possibly CsNO_3 , are candidate vapor species. Unfortunately, because of the small Cs^+ intensity, it was not possible to detect parent ion signals, which can be expected to be weaker than for the Cs^+ fragment. The most probable volatile species is considered to be CsCl because of the likelihood of alkali exchange reactions (such as 3.11) taking place in the mixture. The simultaneous observation of LiCl^+ , NaCl^+ , Na^+ and Cs^+ , as shown in figure 9, is indicative of the presence of such reactions. Near 650°C, other alkali metal ions and halides, including Na_2Cl^+ , KCl^+ , K^+ , and Li^+ , were observed in addition to LiCl^+ , NaCl^+ and Na^+ . During this stage SO_2^+ and its fragments SO^+ and S^+ were also noted.

The appearance of Na_2Cl^+ indicates $(\text{NaCl})_2$ as a vapor species which, in turn, is evidence of a reasonably high NaCl activity in the condensed phase. In fact, at 800°C the NaCl partial pressure, when compared with the pure salt value [JANAF, 1971], indicates an activity of 3×10^{-3} . This value is

about three times the ideal-solution value for the composition given in table 1 which is further evidence of salt-oxide phase separation. Likewise, the CsCl activity is observed to be at least an order-of-magnitude greater than that expected for a homogeneous phase with the composition of table 1. On the other hand, the observed LiCl activity, using the partial pressure data of figure 9 and the JANAF [1971] reference state data, is about 2×10^{-3} . This is much less than expected from the total amount of Li present (see table 1) and indicates only a relatively small conversion of Li_2O to LiCl. The absence of $(\text{LiCl})_2$ in the vapor is also consistent with a low LiCl activity.

Above 800°C , mass spectral ions corresponding to LiBO_2 , NaBO_2 , and O_2 were observed, as shown in figure 10, in addition to the various alkali halide and sulfate-derived ions already noted. In this temperature region, Cs^+ was not observed even though the concentration of CsBO_2 in the glass phase should be comparable to that of the synthetic nuclear waste glass studied by us previously. Thus, it might be inferred that the Cs component of the sludge-frit mixture was lost as the volatile chloride before incorporation into the glass phase.

This conclusion is also supported by the magnitude of the Cs^+ (CsCl) signal in figure 9, which is sufficiently high to cause sample depletion at temperatures less than 800°C . The magnitude of the K^+ ion intensities in figure 10 are also consistent with KCl as the precursor, as formed by the halogen exchange process. At the highest temperatures, KBO_2 is also a possible mass spectral precursor for K^+ . The negative curvature shown in the high temperature range of figure 10, corresponds to a decreasing partial pressure caused by decreased concentration of the various volatile components in the melt.

3.7.2 Knudsen Gravimetric Analysis

A sample of the sludge-frit mixture was subjected to Knudsen gravimetric analysis. These measurements are complimentary to the mass spectrometric analysis since, in principle, the rate of vaporization can be related to the partial pressure of each vapor species. Figure 11 shows the rate of vaporization data and a hypothetical partial pressure which would be valid if only a single species of mass 50 amu was present in the vapor phase. In the present case, the sludge-frit system is complicated by the presence of many significant vapor species. A more detailed analysis can be made by apportioning the total mass loss between the various vapor species observed mass spectrometrically. When this is done, the average vapor molecular weight increases to about 63 amu and the pressure scale of figure 11 is reduced by 10 percent, which is a negligible correction for the present purpose.

3.7.3 TMS Observations

Observations of the sludge-frit mixture, using the TMS system, were carried out at temperatures up to 1400°C using argon carrier gas pressures in the range 0.2 - 0.4 atm. Representative data are given in figure 12. In contrast with the KMS experiments, the high ambient pressure allowed a more uniform release of SO₂ from the melt and troublesome bubble formation was avoided. However, the lower TMS sensitivity precluded observation of low intensity species such as Cs⁺, NaCl⁺, and LiCl⁺.

The Li^+ and Na^+ intensities, shown in figure 12, are derived from several species including LiBO_2 and NaBO_2 as the dominant precursors at the higher temperatures. Near 1150 °C, the parent ions LiBO_2^+ and NaBO_2^+ were also observed. The approximate partial pressures of LiBO_2 and NaBO_2 , given in figure 12, agree reasonably well with our earlier results on a similar glass composition [Hastie et al., 1983], particularly at the higher temperatures. At lower temperatures, the pressure scale is probably a factor of three too high, as may be seen by comparison with figure 11.

Sulfur dioxide was evolved over most of the temperature range but its partial pressure decreased with increasing temperature due to surface depletion between 1050 and 1250 °C before increasing again. The oxygen species was readily detected above 1050 °C but at lower temperatures the presence of S^+ from SO_2 at the same amu interfered with its measurement.

4. Conclusions

The results of the present study indicate a variety of modes by which Cs can be lost to the vapor phase during the nuclear waste glass-forming process. These processes are summarized in simplified form in table 3.

At 190 °C, pure CsO_2CH converts to the oxalate with vaporization loss of about 1 percent of the total Cs.

At 250 °C, under conditions where a pure CsOH condensed phase is present, significant vaporization occurs as CsOH and $(\text{CsOH})_2$.

At 500 °C, with relatively high heating rates, transformation of cesium formate to the oxalate and oxalate to carbonate can occur simultaneously. Under these conditions, Cs vapor transport occurs and appears to be enhanced

by the presence of $\text{H}_2(\text{g})$ or $\text{C}(\text{s})$ produced during the formate-oxalate decomposition. This conclusion is borne out by the vapor pressure measurements on formate-derived carbonate where the data fall in the range expected for reduction of Cs_2CO_3 by C , as calculated from thermochemical data.

At 550 °C, Cs vapor transport can occur, probably as $\text{CsCl}(\text{g})$, by a condensed phase halogen exchange process involving NaCl , CrCl_3 , or CoCl_2 .

At 700 °C, another mechanism whereby the formate-derived carbonate is converted to volatile species occurs by reaction with water vapor to produce $\text{CsOH}(\text{g})$.

At 800 °C, for conditions where condensed Cs_2CO_3 is present, the Cs decomposition pressure is significant.

At 1000 °C, the results of an earlier study [Hastie, et al., 1983] indicate Cs vapor transport in the form of $\text{CsBO}_2(\text{g})$. However, in the presence of halides most of the Cs can be lost at lower temperatures, as indicated above. It appears, therefore, that the presence of halides is detrimental to the stable incorporation of cesium in the borosilicate glass.

The degree of cesium vapor transport and the relative significance of the various reactions will, in practice, depend on the temperature-time history of the vitrification process. Development of thermodynamic mixing models, based on the present experimental data, and a solution model [eg. see Hastie and Bonnell, 1985] is required to model the plant-scale process.

Table 3

Summary of processes leading to Cs vapor transport

Material	Temperature °C	Process
cesium formate	190	$\text{CsO}_2\text{CH(s)} \longrightarrow \text{Cs(g)} \text{ and/or } \text{CsO}_2\text{CH(g)}$
cesium hydroxide	250	$\text{CsOH(c)} \longrightarrow \text{CsOH(g)}$
cesium formate	500	$\text{Cs}_2\text{CO}_3\text{(c)} + \text{C(s)} \longrightarrow \text{Cs(g)}$
sludge-frit	550	$\text{Cs}_2\text{O(c)} + \text{NaCl(c)} \longrightarrow \text{CsCl(g)}$
sludge-frit	700	$\text{Cs}_2\text{CO}_3\text{(s)} + \text{H}_2\text{O} \longrightarrow \text{CsOH(g)}$
cesium carbonate	800	$\text{Cs}_2\text{CO}_3\text{(l)} \longrightarrow \text{Cs(g)}$
sludge-frit	1000	$\text{CsBO}_2\text{(c)} \longrightarrow \text{CsBO}_2\text{(g)}$

5. Acknowledgments

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Reisman, A. (1958). J. Phys. Chem., 80, 3558.

Schoonmaker, R.C. and Porter, R.F. (1959). J. Chem. Phys., 31, 830.

7. Figure Captions

1. Comparison of KGA results (open squares, and least squares curve) for apparent $P(\text{CsOH})$ over $\text{CsOH}(\text{s},\text{l})$ with the JANAF [1971] monomer and dimer data, as labelled.
2. Ratio of $\text{I}^+\text{Cs}/\text{I}^+\text{O}_2$ and $\text{I}^+\text{Cs}/\text{I}^+\text{Cs}_2\text{O}$ mass spectral ion intensity data from Gusarov et al. [1967a] using an initial condensed phase with composition $\text{Cs}_2\text{O}_2(\text{c})$. I^+ is mass spectral ion-intensity in arbitrary units.
3. Comparison of KMS and literature (Lamoreaux and Hildenbrand, 1984) equilibrium constant data for decomposition of $\text{Cs}_2\text{O}(\text{g})$ to $\text{Cs}(\text{g})$ and $\text{O}_2(\text{g})$; data obtained from evaporation of $\text{Cs}_2\text{O}(\text{l})$ with increasing and decreasing temperature cycles.
4. KGA Cs partial-pressure data, and fitted curves, obtained over $\text{Cs}_2\text{CO}_3(\text{c},\text{l})$ (triangles) and formate-derived Cs_2CO_3 (squares) compared with literature data over Cs_2CO_3 (Glushko et al. [1978-1982]) and the calculated $\text{Cs}_2\text{CO}_3 + \text{C}$ reaction pressure curve.
5. KMS Cs partial-pressure data, and fitted curve, obtained over $\text{Cs}_2\text{CO}_3(\text{c},\text{l})$ (square symbols). Literature curve is from Glushko et al. [1978-1982].
6. Ion-intensity data for the reaction of $\text{Cs}_2\text{CO}_3(\text{c})$ with H_2O to form CsOH at 1010 K (see reaction (3.10) in text).
7. Comparison of experimental and calculated equilibrium constant data for reaction (3.10) (see text) involving $\text{Cs}_2\text{CO}_3(\text{c})$ and $\text{H}_2\text{O}(\text{g})$.
8. TMS ion-intensity and approximate partial-pressure data for decomposition of formate-related constituents from sludge.
9. KMS mass spectral ion-intensity and approximate partial-pressure data for relatively low temperature vapor species over the sludge-frit mixture. The partial-pressure scale relates to the species given in parentheses.
10. KMS mass spectral ion-intensity, and approximate partial-pressure data for high temperature vapor species over the sludge-frit mixture. The partial-pressure scale relates to the species given in parentheses.
11. KGA rate of vaporization and apparent vapor pressure data obtained over the sludge-frit mixture for a hypothetical vapor molecular weight of 50 amu.
12. Selected TMS mass spectral ion-intensity and approximate partial-pressure data for high temperature vapor species over the sludge-frit mixture. The partial-pressure scale relates to the species given in parentheses.

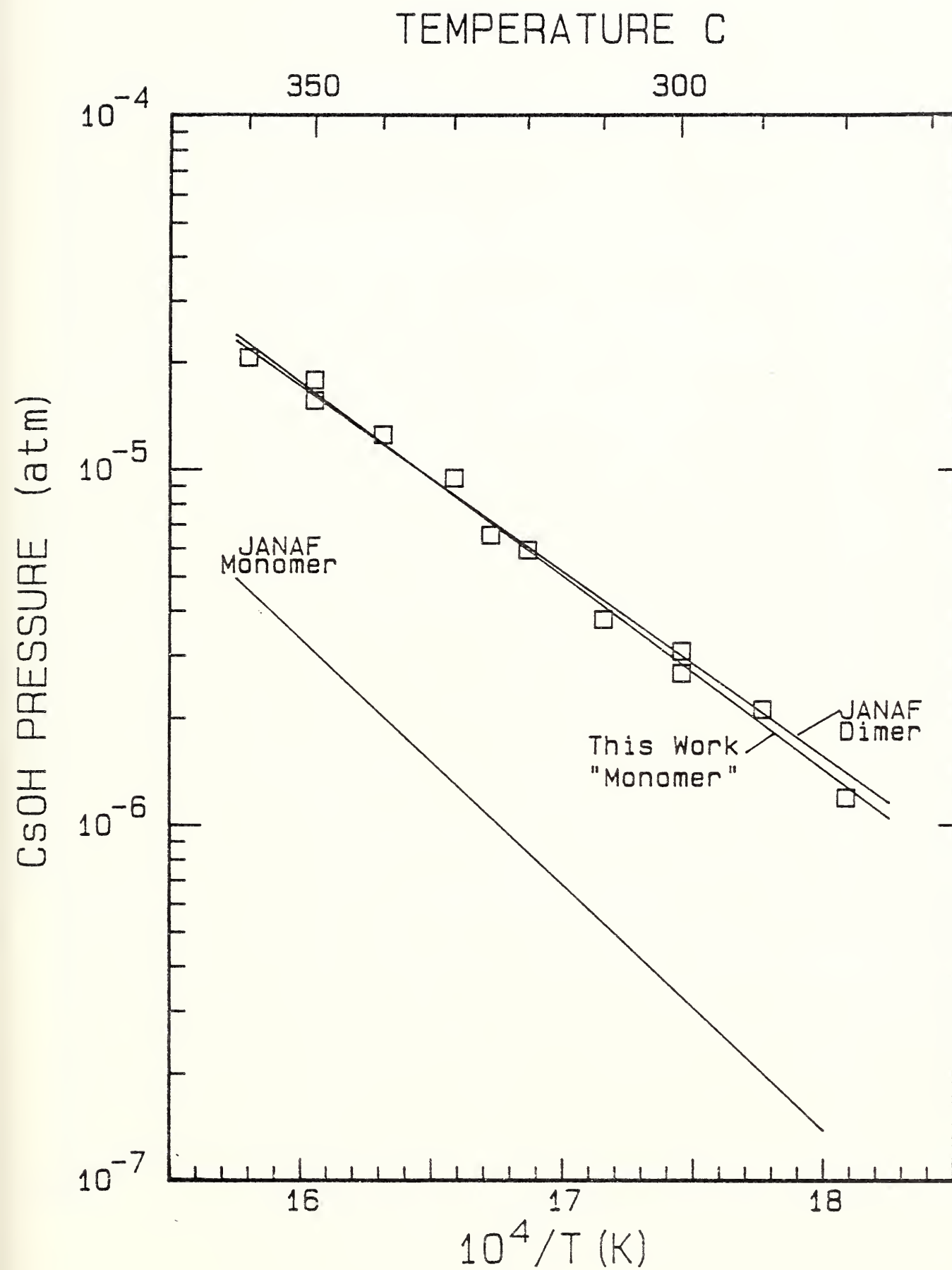


Figure 1

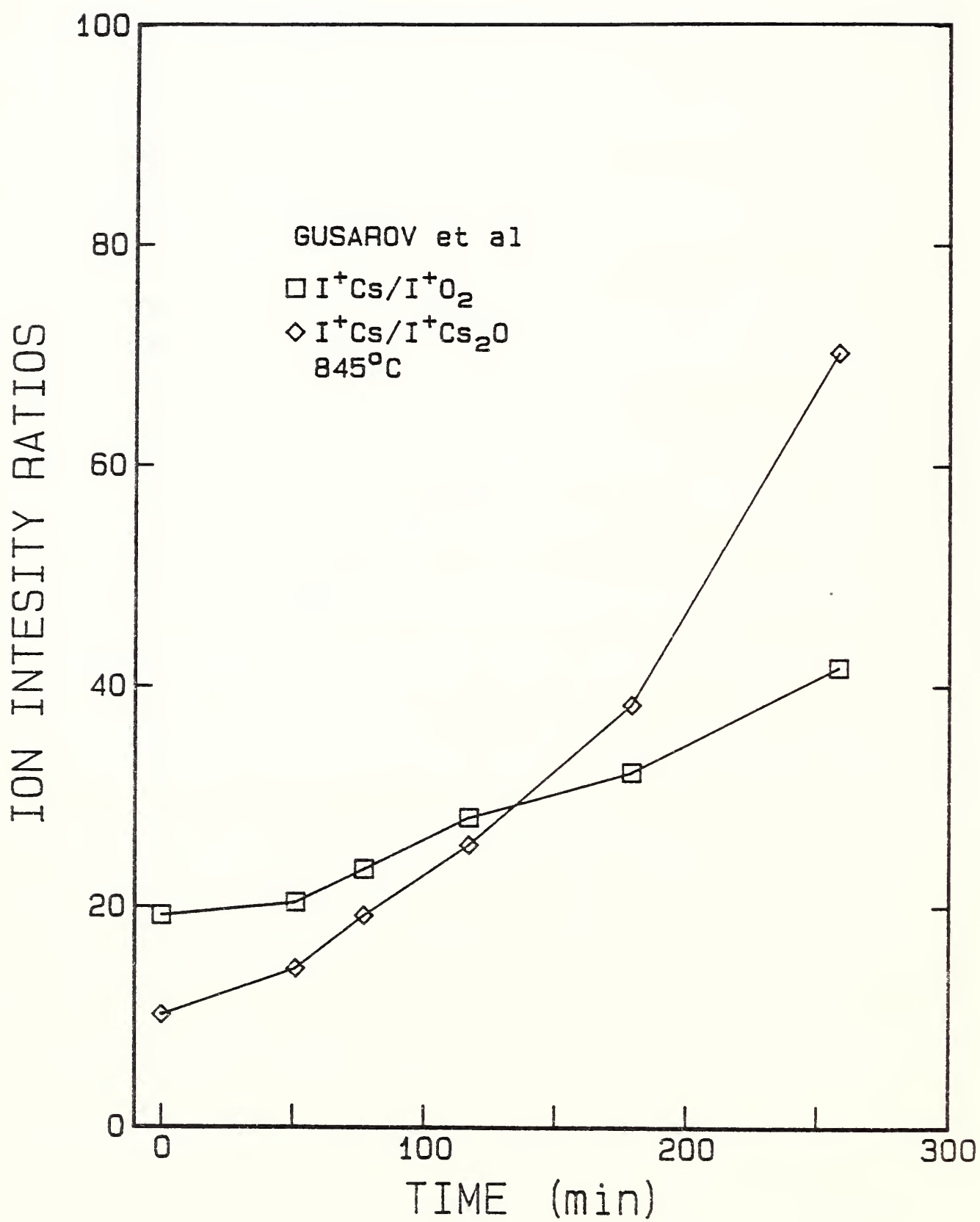


Figure 2

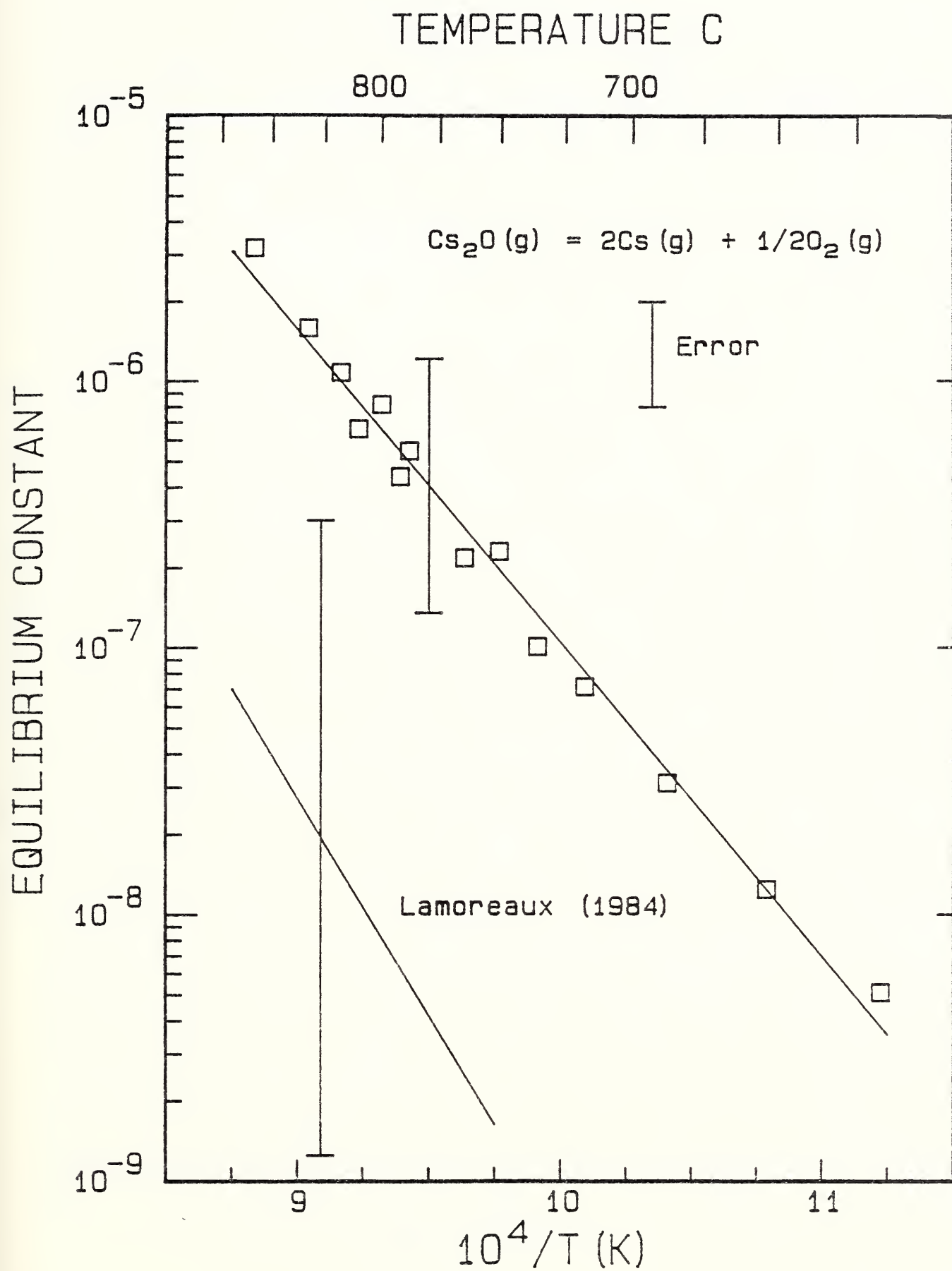


Figure 3

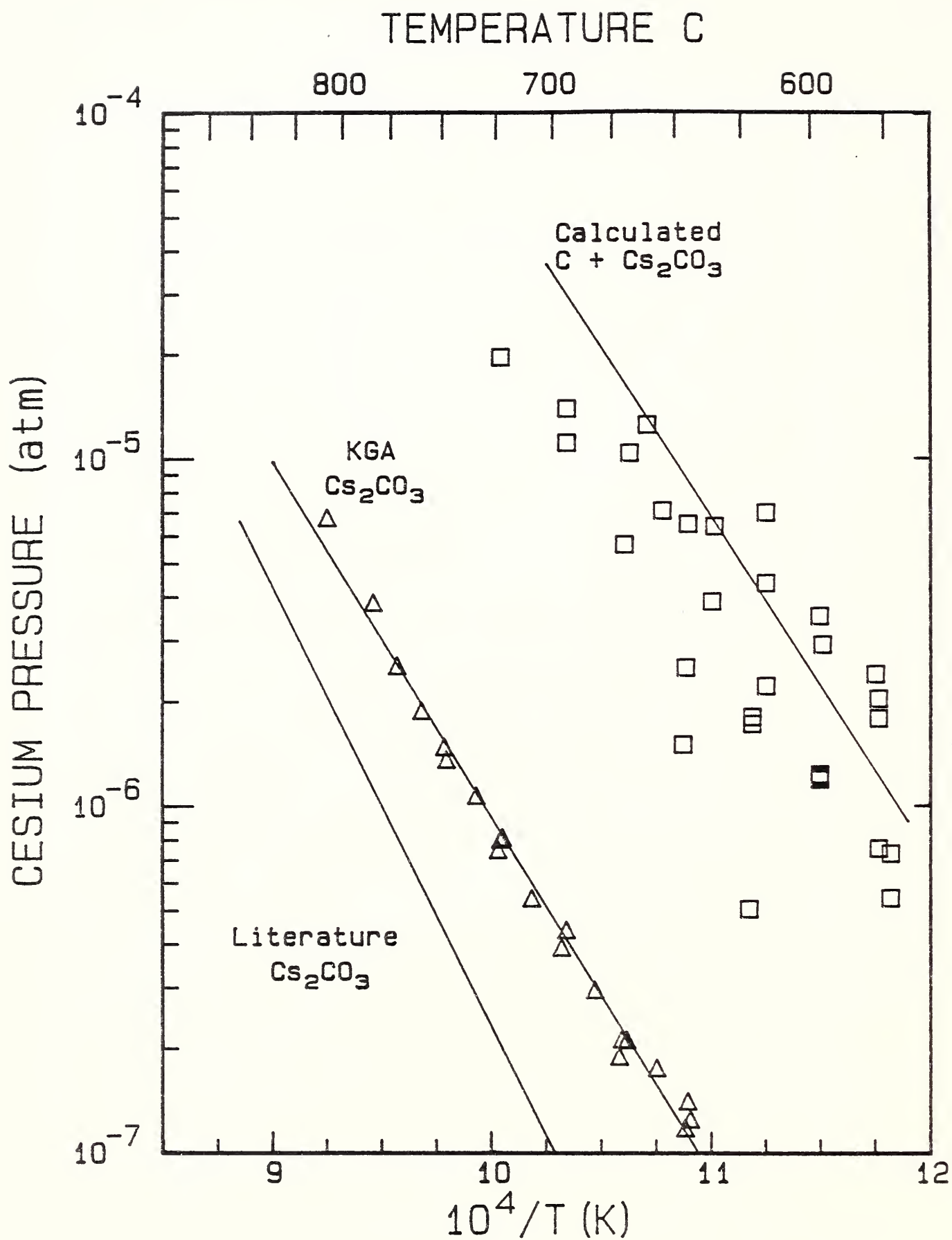


Figure 4

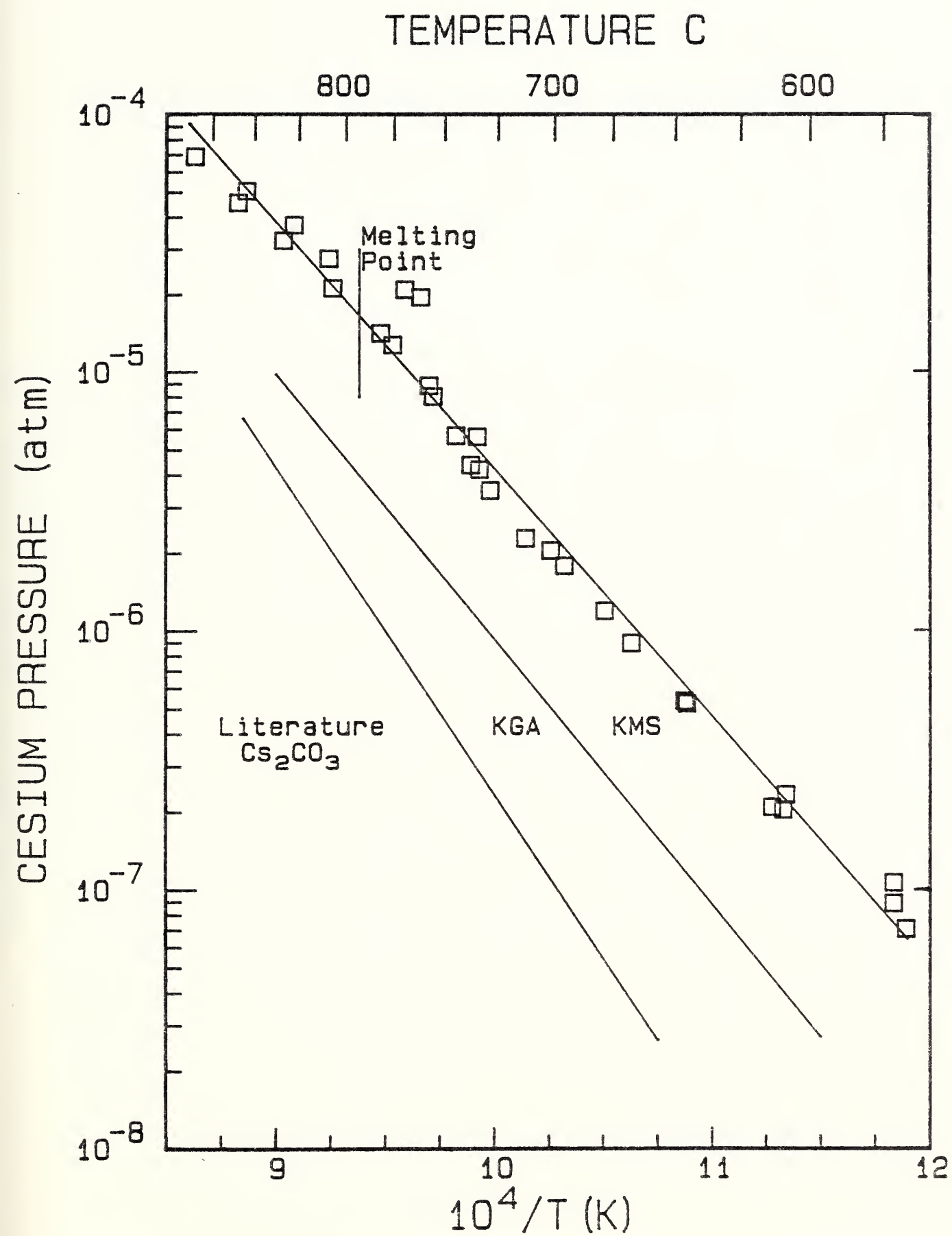


Figure 5

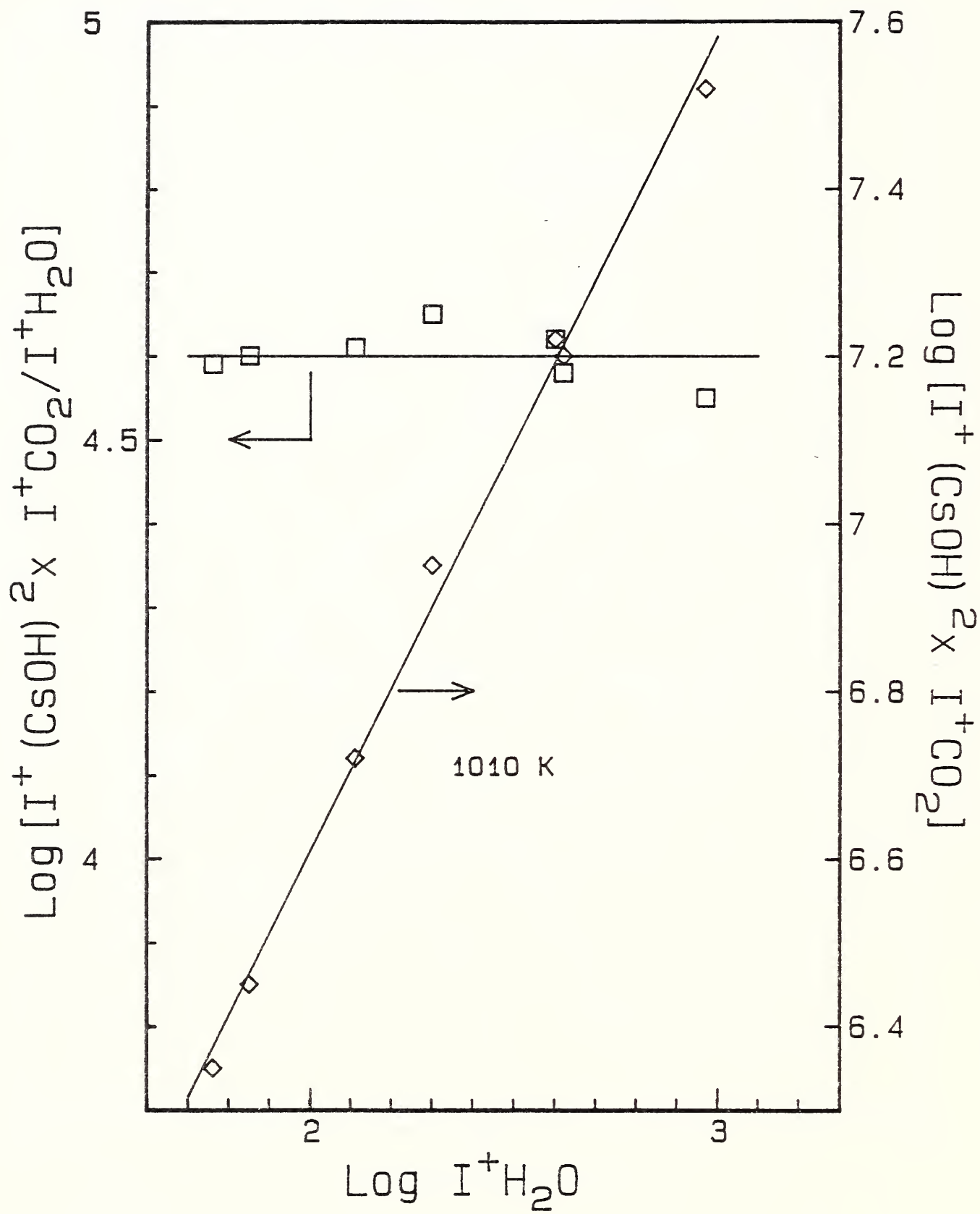


Figure 6

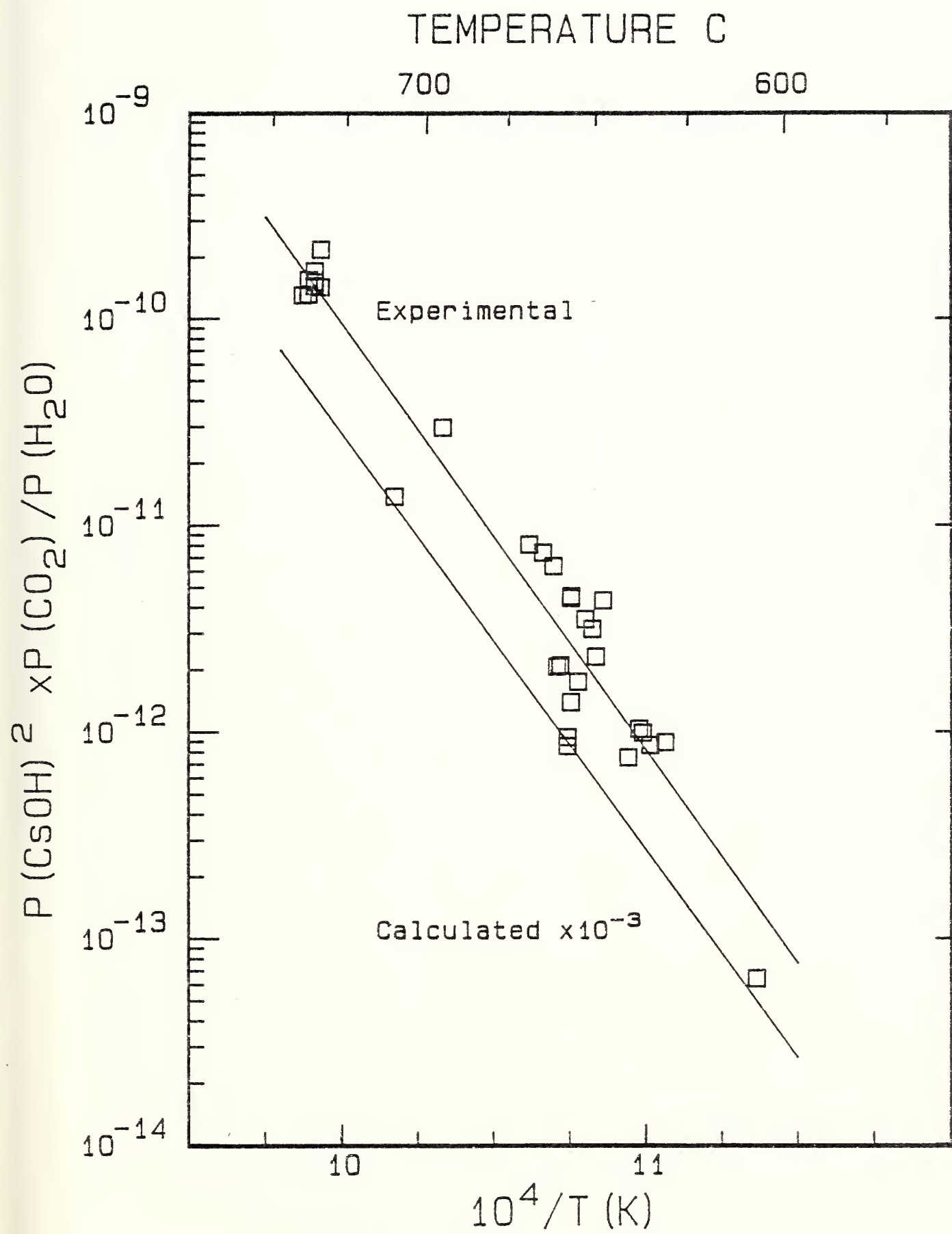


Figure 7

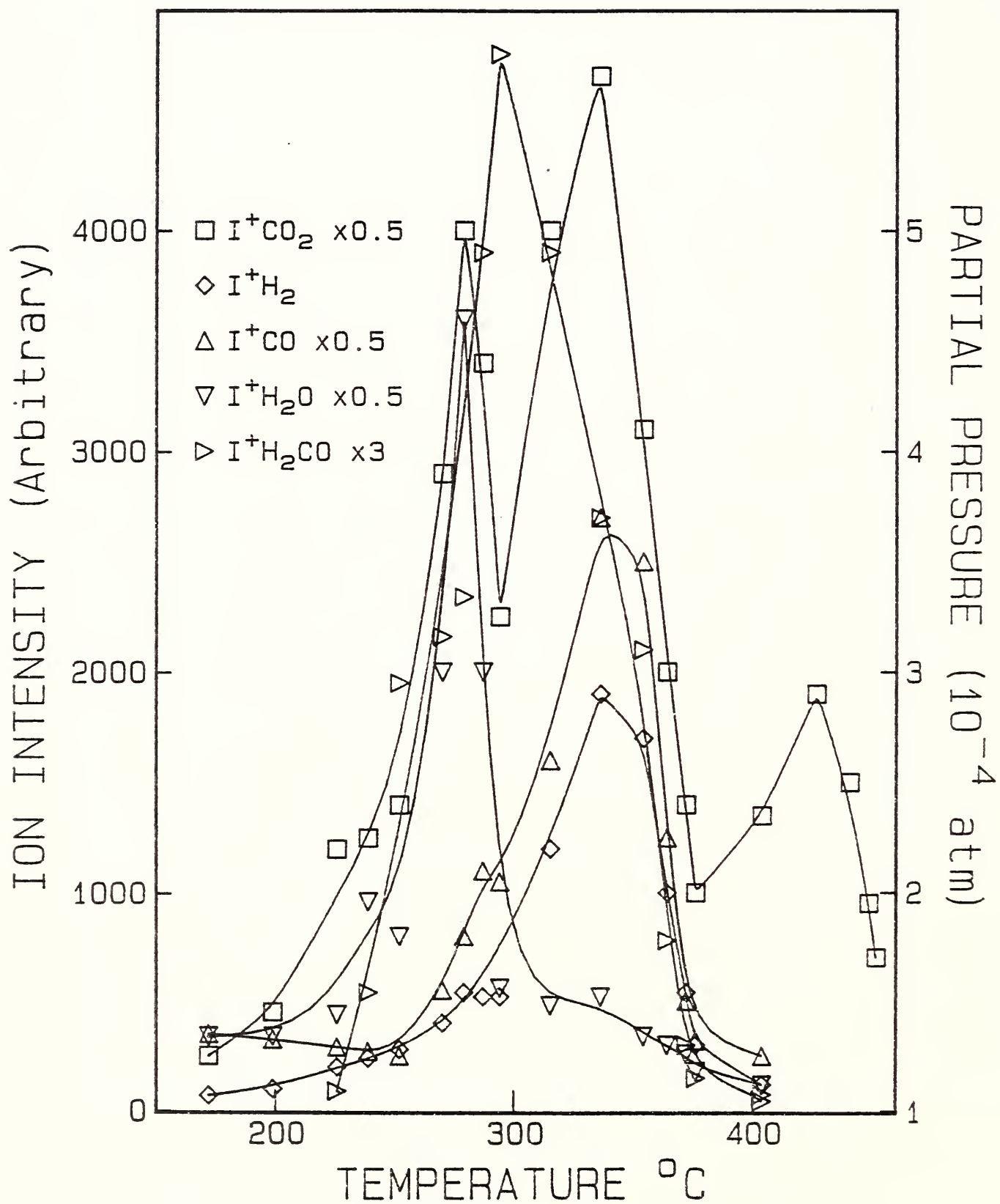


Figure 8

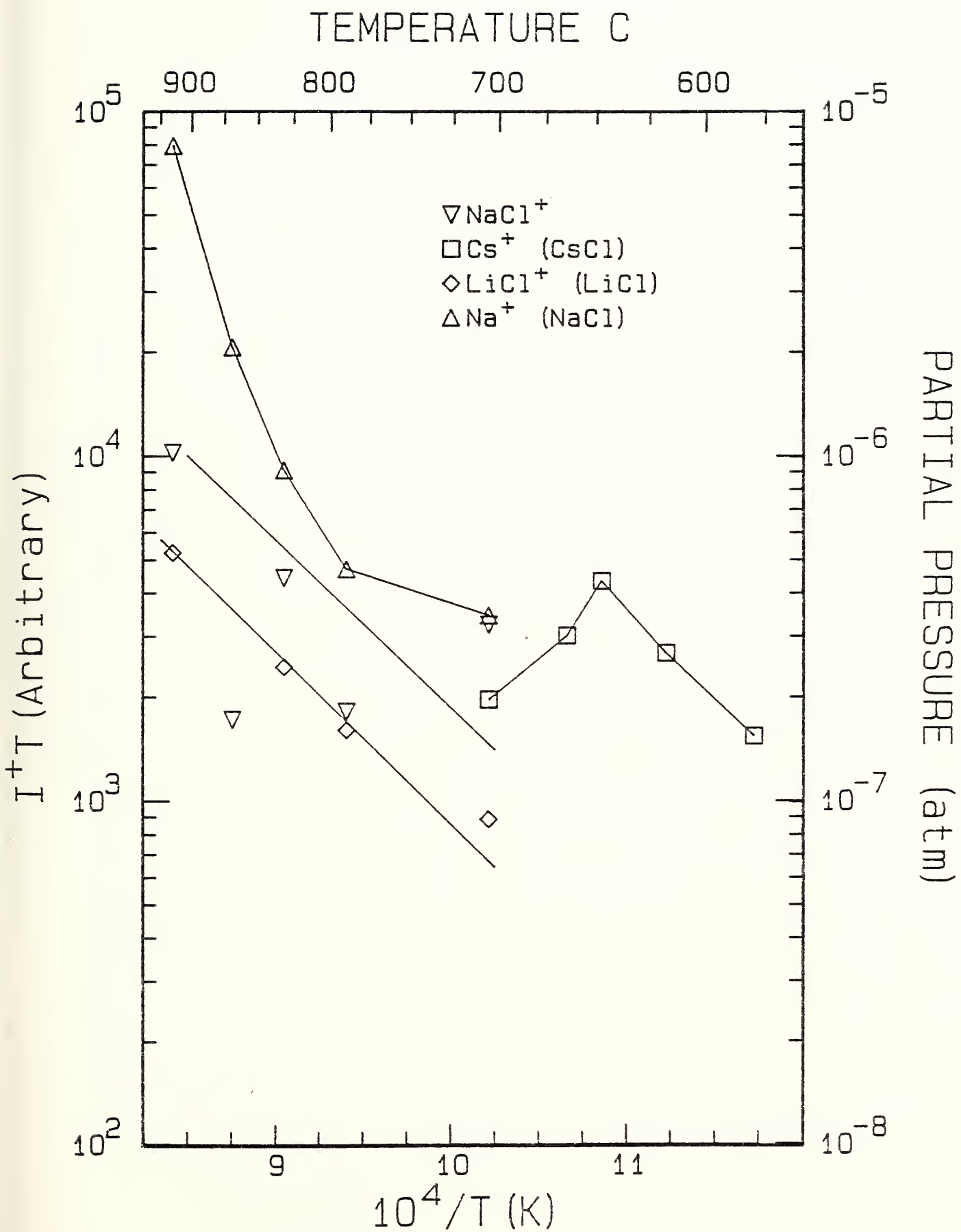


Figure 9

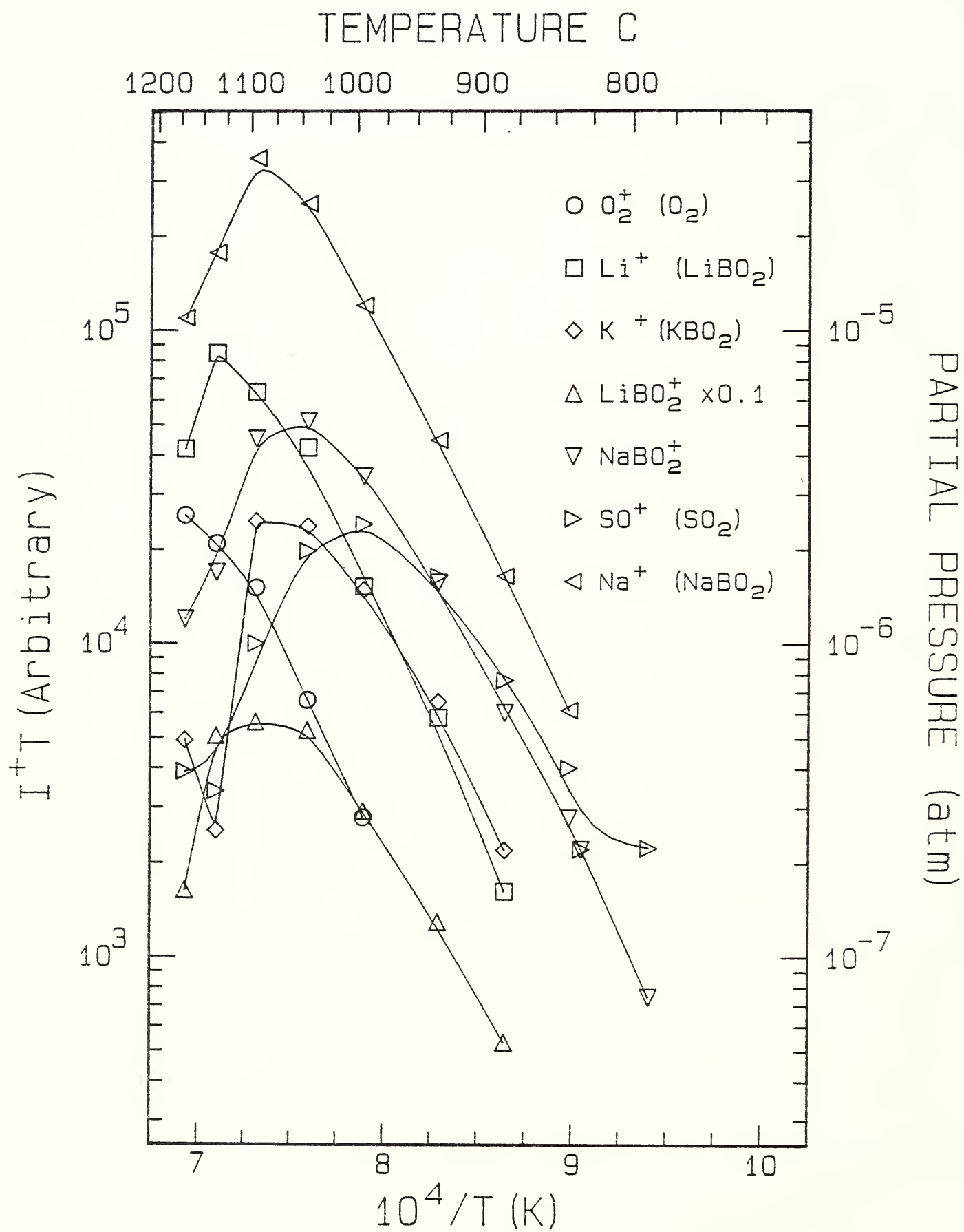


Figure 10

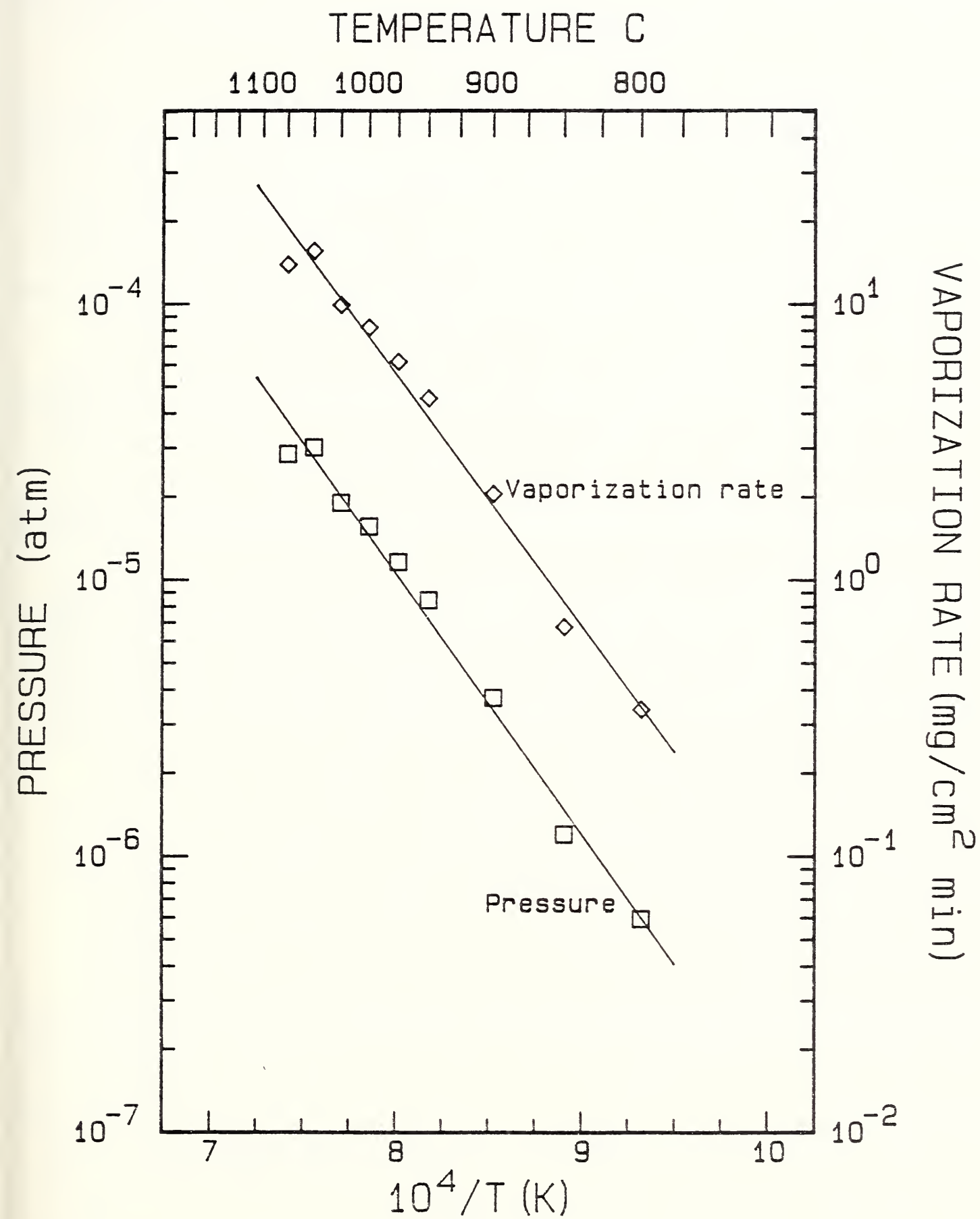


Figure 11

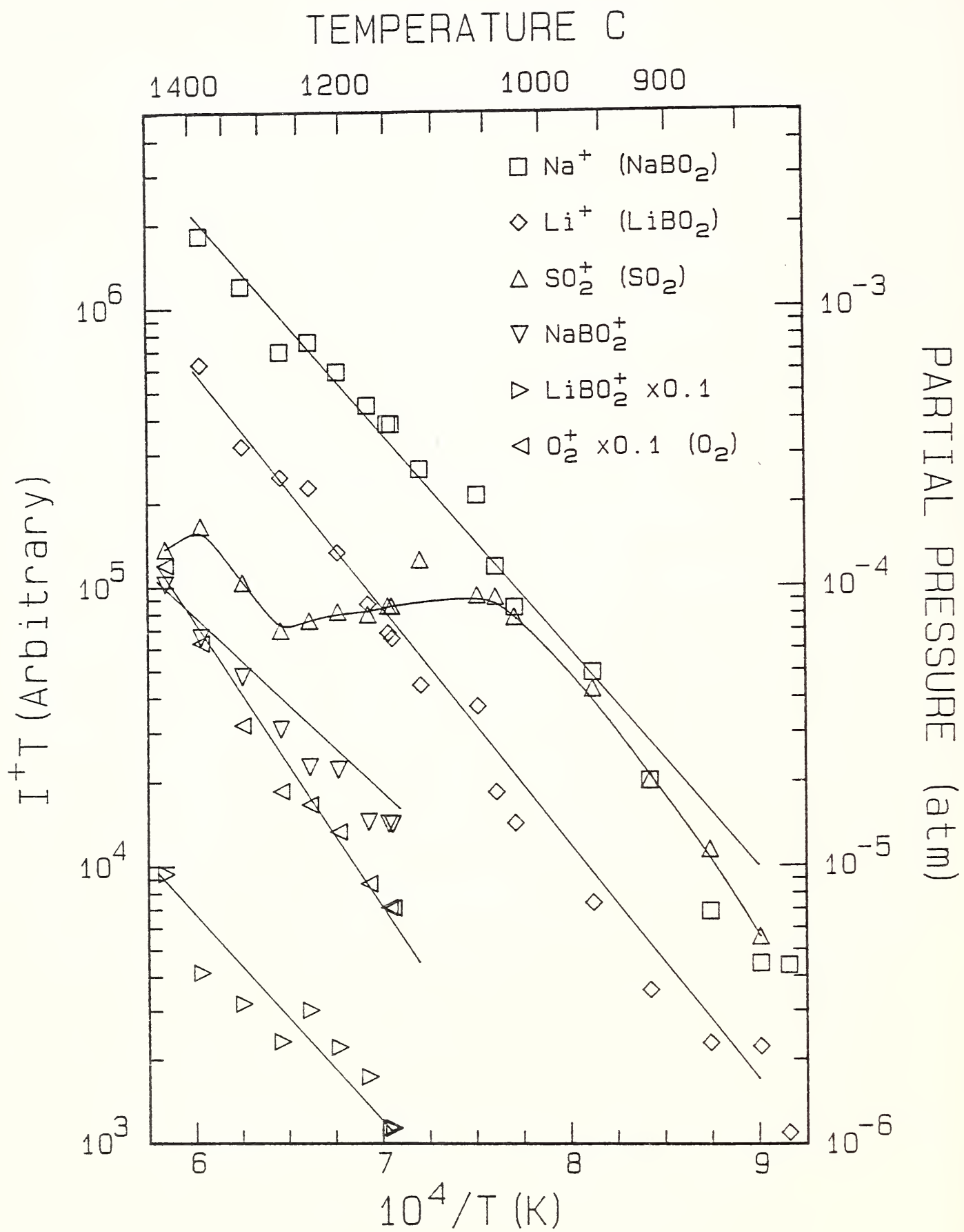


Figure 12

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5. AUTHOR(S) E.R. Plante and J.W. Hastie			
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11. ABSTRACT (A 200-word or less factual summary of most significant information. If document includes a significant bibliography or literature survey, mention it here) Undesirable losses of radionuclides occur by vaporization during processing of nuclear waste glass. This study examines the phase and chemical transformations that occur during the initial steps of waste incorporation into a borosilicate glass host. Specialized mass spectrometric and thermogravimetric techniques were used to monitor the vapor transport species over a range of temperature and composition. Significant vapor transport of Cs was found and its enhancement by the presence of carbon, halogen, and water vapor was also noted.			
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